

## Extrinsic S/C Prop.: Free Carrier Statistics

### Free Carrier Statistics:

- ✓ Previously, we determined the eqns:
  - $n = f(T)$
  - $p = f(T)$
- ✓ Carrier concentrations as a function of temperature in intrinsic S/C's.
- ✓ We found that:

$$n_i^2 = N_c N_v e^{-E_g/kT}$$

- ✓ Since we have been discussing impurities in S/Cs (extrinsic S/Cs), we would like to find  $n$  and  $p$  for the extrinsic case.
- ✓ That is, we want to determine how many free  $e^-$ 's and  $h^+$ 's exist near their respective band edges as a function of:
  - Impurity concentration
  - Temperature
- ✓ Called "Free Carrier Statistics", this is again performed using:
  - Statistical Mechanics
  - Quantum Mechanics

## Extrinsic S/C Prop.: Free Carrier Statistics

### Let's define our extrinsic S/C system:

- ✓ One type of donor (when I say "donor" or "acceptor", I mean shallow-level donor or acceptor):
  - The donor has energy,  $E_D$
  - The donor concentration,  $N_D$
- ✓ One type of acceptor
  - The acceptor has energy,  $E_A$
  - The acceptor concentration,  $N_A$
- ✓ The position of the Fermi energy level,  $E_f$ , will be determined self-consistently by the:
  - Temperature,  $T$
  - Impurity concentrations,  $N_D$  &  $N_A$
  - Band gap energy,  $E_g$
- ✓ One additional factor to introduce: Degeneracy,  $g$ , in  $k$ -space:
  - Donors:
    - $g = 2$ , (spin  $1/2$ )  $m_j = \pm 1/2$  (2 spin states)
  - Acceptors:
    - $g = 4$ , (spin  $3/2$ )  $m_j = \pm 3/2$  (4 spin states)
      - »  $m_j$  is the total angular momentum

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- ❑ We now can ask: “What is the probability that a donor level at energy  $E_D$  with degeneracy  $g$  is occupied?”
- ❑ In order to answer this question, we need to use a distribution statistics of  $e^-$ 's &  $h^+$ 's
- ❑ What distribution function would you suggest?

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- ❑ The Fermi-Dirac distribution function is (w/degeneracy  $g$ ):

$$F(E) = \frac{1}{1 + g e^{(E-E_f)/kT}}$$

- ❑ The probability for occupancy of a donor level at energy  $E_D$  is:

$$F(E_D) = \frac{1}{1 + g_D e^{(E_D-E_f)/kT}}$$

- ❑ The fraction of donor levels,  $N^*$ , which are occupied by  $e^-$ s is:

$$N^* = N_D \cdot F(E_D) = \frac{N_D}{1 + g_D e^{(E_D-E_f)/kT}}$$

- ❑ Let's choose a specific case in which:

$$N_D > N_A$$

- ❑ What “type” of extrinsic S/C do we have?

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- ❑ The maximum number of e<sup>-</sup>'s that can be promoted to the CB is:

$$N = N_D - N_A$$

- ❑ There are  $N_A$  e<sup>-</sup>'s compensated by acceptors.
- ❑ Since the acceptor states lie very close to the VB, the compensated e<sup>-</sup>'s in the acceptor states need nearly the  $E_g$  to be promoted to the CB.
- ❑ Therefore, they can be safely neglected at low temperatures:  $T < \text{room temperature}$
- ❑ Let:  $n$  = number or concentration of e<sup>-</sup>'s in CB
- ❑ Thus:  $N = N_D - N_A = N^* + n$

✓ Where  $N^*$  is concentration of donor states occupied by e<sup>-</sup>'s

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- ❑ Since:  $N^* = N_D \cdot F(E_D)$
- ❑ Substituting  $N^*$  in the previous equation:

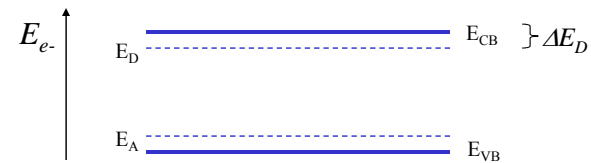
$$N = \frac{N_D}{1 + g_D e^{(E_D - E_f)/kT}} + n$$

$$= \frac{N_D}{1 + g_D e^{E_D/kT} e^{-E_f/kT}} + n$$

- ❑ Thus:

$$\frac{N - n}{N_D} = \frac{1}{1 + g_D e^{E_D/kT} e^{-E_f/kT}}$$

- ❑ Let's draw the energy band diagram when the S/C is compensated and define certain energy levels.



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- Continuing with our previous equation and taking the reciprocal of both sides, we have:

$$\frac{N_D}{N-n} = 1 + g_D e^{E_D/kT} e^{-E_f/kT}$$

$$\frac{N_D}{N-n} - 1 = g_D e^{E_D/kT} e^{-E_f/kT}$$

$$\frac{N_D - N + n}{N-n} = g_D e^{E_D/kT} e^{-E_f/kT} \quad (i)$$

- From our earlier determination of the number of e<sup>-</sup>'s in the CB, we found that:

$$n = N_C e^{-(E_c - E_f)/kT}$$

- Where we assumed the *Boltzmann Approximation*.
- Thus, we are assuming that the S/C is nondegenerate.

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- Solving for E<sub>f</sub> we have:

$$E_f = E_c + kT \ln \left( \frac{n}{N_C} \right)$$

- Substituting E<sub>f</sub> into equation. (i), we obtain:

$$\frac{N_D - N + n}{N-n} = g_D e^{-(E_c - E_D)/kT} \frac{N_C}{n}$$

Let:  $\Delta E_D = E_c - E_D$  &  $\Theta = N_C g_D e^{-\Delta E_D/kT}$

So:  $\frac{n(N_D - N + n)}{N-n} = \Theta$  (ii)

- With some algebra, we obtain a quadratic equation in n:

$$n^2 + (N_D + \Theta - N)n - N\Theta = 0$$

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- ❑ To determine the number of e<sup>-</sup>'s in CB, solve for n.
- ❑ Using the Quadratic equation, solve for n:

$$n = \frac{1}{2} \left\{ -(N_D + \Theta - N) \pm \sqrt{(N_D + \Theta - N)^2 + 4\Theta N} \right\}$$

$$= \frac{1}{2} \left\{ -(N_D + \Theta - N) \pm (N_D + \Theta - N) \sqrt{1 + \frac{4\Theta N}{(N_D + \Theta - N)^2}} \right\}$$

$$= \frac{1}{2} (N_D + \Theta - N) \left\{ \pm \sqrt{1 + \frac{4\Theta N}{(N_D + \Theta - N)^2}} - 1 \right\}$$

- ❑ Since  $n \geq 0$ , we take only the positive solution.

$$n = \frac{1}{2} (N_D + \Theta - N) \left\{ \sqrt{1 + \frac{4\Theta N}{(N_D + \Theta - N)^2}} - 1 \right\}$$

Where:  $\Theta = N_C g_D e^{-\Delta E_D / kT}$

- ❑ This last expression is good to within several  $kT$  of the CB b/c the Boltzmann Approximation was used.

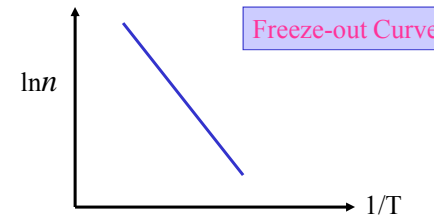
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We can use equation (ii) to examine several Cases:

- ❑ Case 1:  $N_A = 0$  therefore  $N = N_D$ .
  - ✓ @ low T,  $n \ll N_D$ .
  - ✓  $n$  = concentration of CB e<sup>-</sup>'s, a.k.a. the free carrier concentration

$$n \approx \sqrt{g_D N_D N_C} e^{-\Delta E_D / kT} = (g_D N_D N_C)^{1/2} e^{-\Delta E_D / 2kT}$$

✓ Plot of ln(n) vs 1/T

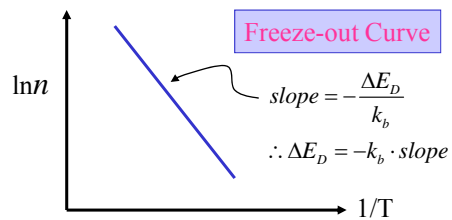


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- Case 2:  $N_A \neq 0$ .  $T \sim$  low, thus  $n \ll N_A < N_D$

$$n \approx g_D N_C \frac{N_D - N_A}{N_A} e^{-\Delta E_D / kT}$$

- Plot of  $\ln(n)$  vs  $1/T$



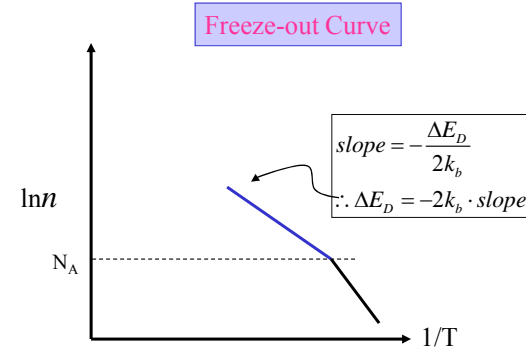
- This region is also known as the “Full-slope regime”.

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- Case 3: Increase  $T$  such that  $N_A \ll n < N_D$

$$n \approx (g_D N_C N_D)^{1/2} e^{-\Delta E_D / 2kT}$$

- Similar to case 1, we find that the slope is  $E_D/2k$  or “half-slope”.
- Hence, the region in blue below is known as the “half-slope regime”.

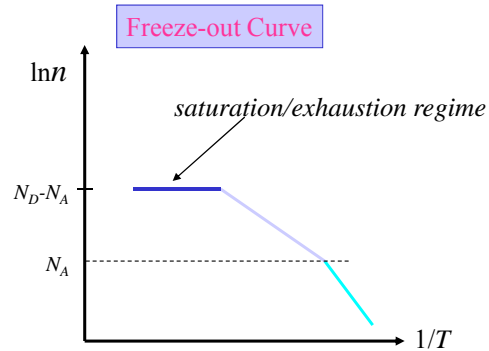


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Case 4:

- ✓ As  $T$  continues to increase, a region is attained in which the free  $e^-$  concentration,  $n$ , remains constant.
- ✓ This is known as the
  - o Exhaustion regime
  - o Saturation regime
- ✓ In this condition,  $n$  is given by:

$$n = N_D - N_A$$



- ✓ **Question:** What happens at higher temps beyond the saturation regime?

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Fermi Energy Level of a Doped S/C as a Function of Temperature

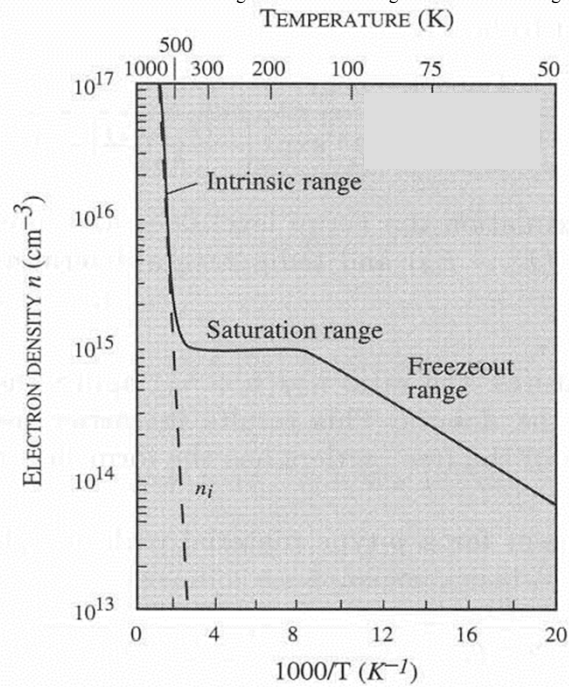
- ✓ Assume equilibrium
- ✓ At temperatures sufficiently high for the S/C to be intrinsic,  $E_f$  lies near mid-gap.
- ✓ As  $T$  decreases  $E_f$  moves either toward the:
  - o CB for n-type
  - o VB for p-type
- ✓  $E_f$  coincides with the majority dopant energy level at the  $T$  for which 50% of the majority dopants are frozen out (i.e., neutral).

$$\begin{aligned} \therefore N_D &= N_D^o - N_D^+ \\ N_D^o &= N_D \end{aligned}$$

- ✓ For *uncompensated* S/C (ideal case),  $E_f$  would approach a position close to the mid-point between the dopant level and a band edge (CB or VB).
- ✓ For compensated S/Cs, a fraction of the majority dopants remain ionized at the lowest  $T$  and  $E_f$  will be very close to the majority dopant energy level.

□ In-class Exercise:

- ✓ Determine if n- or p-type?
- ✓  $|N_D - N_A|$ ?
- ✓ Minority dopant concentration?
- ✓ Full slope or half slope?
- ✓  $E_g$ ? What is the S/C material?
- ✓  $E_D$  or  $E_A$ ? What is the dopant?
- ✓ Show where intrinsic region and extrinsic region and saturation regimes are.



□ Freezeout Curve

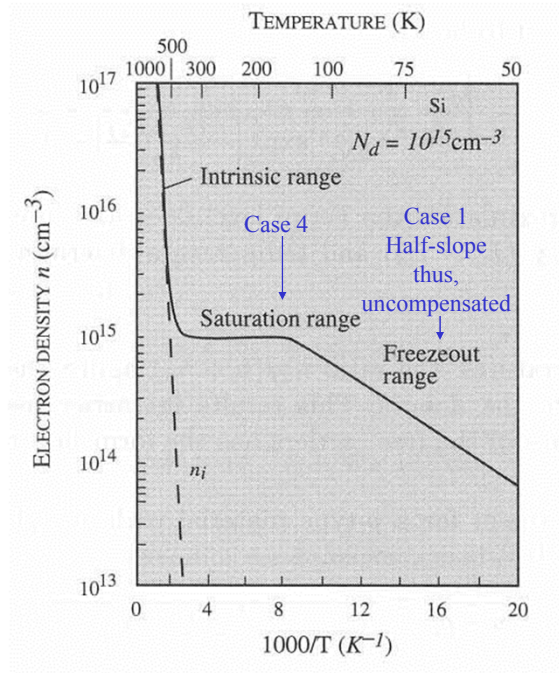


Figure 2.15: Electron density as a function of temperature for a Si sample with donor impurity concentration of  $10^{15}$  cm<sup>-3</sup>. It is preferable to operate devices in the saturation region where the free carrier density is approximately equal to the dopant density.



Freezeout Curve

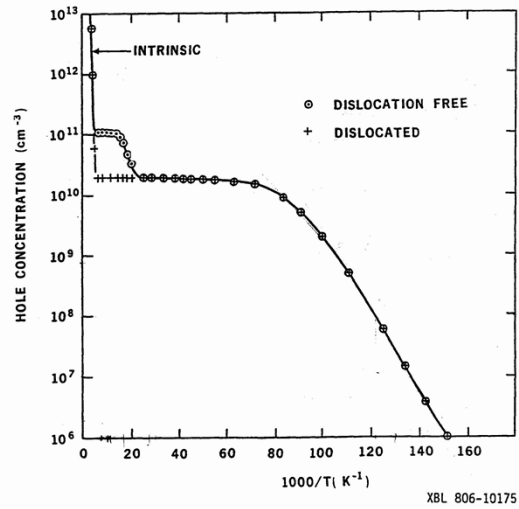


Fig. 21. Arrhenius plots of the free hole concentrations of a dislocation-free and a dislocated piece of single crystal high-purity Ge grown in a H<sub>2</sub> atmosphere.

McCluskey & Haller, *Dopants & Defects in Semiconductors* (CRC Press, 2012) fig. 4.13, p. 120  
After: Haller, Hansen, & Goulding, *Adv. Phys.* (1981)